# A NOVEL METAL CLUSTER IN THE TETRAMETHYLAMMONIUM SALT OF THE DIOCTAHEDRAL TRIACONTACARBONYL DODECARHODATE DIANION, $\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]^{2-}$ 

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SUMMARY
The crystal and molecular structure of the tetramethylammonium salt of the triacontacarbonyldodecarhodate dianion $\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]_{2}\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]$ has been determined from three-dimensional X-Ray data collected by counter methods. The structure has been refined by least squares on $1800 F_{\mathrm{o}}$ 's down to an $R$ factor of 0.05 . This salt crystallizes in the monoclinic space group $P 2_{1} / c\left(C_{2 h}^{5}\right)$ with the following cell parameters: $a=9.491(5) \AA, b=17.185(7) \AA, c=17.756(7) \AA, \beta=96^{\circ} 57^{\prime}\left(15^{\prime}\right)$, $Z=2$.

The structure is composed of anionic and cationic layers. The anion is a centrosymmetric dimer of the $\left[\mathrm{Rh}_{6}(\mathrm{CO})_{15}\right]^{-}$unit. The catenation between the symmetry related sub-units takes place by a metal-metal bond and two CO groups bridged over it. The coordination pattern around the rhodium atoms is very similar to that found in the neutral molecule $\mathrm{Rh}_{6}(\mathrm{CO})_{16}$ with the exception of the catenating Rh atoms which display nine-fold coordination.

The overall symmetry of the anion is $D_{2 h}$. The metal cluster exhibits a range of $\mathrm{Rh}-\mathrm{Rh}$ distances from $2.68 \AA$ to $2.85 \AA$. The independent CO groups fall into 3 classes: 10 are linearly bonded, 4 are triply asymmetrically bridged on alternating octahedral faces and 1 is connecting the two octahedral sub-units through a symmetric double bridge.

## INTRODUCTION

The syntheses and structures of the following carbonylic clusters of cobalt atoms, $\mathrm{Co}_{6}(\mathrm{CO})_{16},\left[\mathrm{Co}_{6}(\mathrm{CO})_{15}\right]^{2-}$ and $\left[\mathrm{Co}_{6}(\mathrm{CO})_{14}\right]^{4-}$ have been described recently ${ }^{1-6}$. On the basis of these results and the existence of $\mathrm{Rh}_{6}(\mathrm{CO})_{16}{ }^{7}$, the existence of corresponding anions was expected for rhodium.

Dichlorotetracarbonyldirhodium, $\mathrm{Rh}_{2}(\mathrm{CO})_{4} \mathrm{Cl}_{2}$, when treated with nucleophilic agents in presence of carbon monoxide, gives the analogues of the aforementioned cobalt clusters, viz. $\mathrm{Rh}_{6}(\mathrm{CO})_{16}$ (ref. 8), $\left[\mathrm{Rh}_{6}(\mathrm{CO})_{14}\right]^{4^{-}}$and probably $\left[\mathrm{Rh}_{6}(\mathrm{CO})_{15}\right]^{2-}$ (ref. 9). The products have been characterized by chemical, spectroscopic and crystallographic examination. In addition to these, other products have been isolated which contain more complicated clusters, for example, $\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]^{2-}$ (ref. 10) and an anion containing seven metal atoms ${ }^{11}$.

Questions about the molecular complexity, stereochemistry and the nature of the bonding in these new anions are of much interest, and so a detailed X-ray structural investigation of both species was undertaken. We shall discuss the structure of the heptarhodate anion in a separate publication, but we now describe the results obtained for tetramethylammonium dodecarhodate $\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]_{2}\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]$.

EXPERIMENTAL
The salt $\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]_{2}\left[\mathrm{Rh}_{12}(\mathrm{CO})_{3 \mathrm{c}}\right]$ (formula weight 2224) crystallizes in violet-black prisms. Preliminary precession photographs indicated a monoclinic symmetry. The systematic absences, $h 0 l$ with $l$ odd and $0 k 0$ with $k$ odd, indicate a space group $P 2_{1} / c\left(N .14, C_{2 h}^{5}\right)$. The unit cell parameters, determined from film data and refined on the PAILRED diffractometer with the $\omega$-lag procedure outlined by Hornstra ${ }^{12}$, have the following values at $20^{\circ}: a=9.491(5) \AA, b=17.185(7) \AA, c=$ $17.756(7) \AA, \beta=96^{\circ} 57^{\prime}\left(15^{\prime}\right), V=2874.8 \AA^{3}$. In these measurements the wavelength of $\mathrm{MoK} \alpha$ radiation was taken as $0.70930 \AA$. The experimental crystal density, 2.50(3) $\mathrm{g} \cdot \mathrm{cm}^{-3}$ measured by a picnometer using toluene as picnometric liquid, is in agreement with the value of $2.56 \mathrm{~g} \cdot \mathrm{~cm}^{-3}$ computed for two formula units per cell. The crystallographically independent unit is then $\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]\left[\mathrm{Rh}_{6}(\mathrm{CO})_{15}\right]$.

The intensity measurements were done on a PAILRED linear equi-inclination diffractometer. The incident beam was $\mathrm{MoK} \alpha$ radiation as obtained from a monochromatizing single crystal ${ }^{1,2}$ of silicon cut on [111]. The crystal used in the data collection was a thin parallelepiped with dimensions of $0.08 \times 0.11 \times 0.40 \mathrm{~mm}$; its direction of maximum elongation, coincident with the $a$ axis of the cell; was oriented along the $\omega$-axis of the instrument in order to minimize absorption effects. The integrated intensities of eight levels of the reciprocal lattice, from $0 k l$ to $7 k l$, were collected within a range $3^{\circ}<r<52^{\circ}$, the diffraction profiles were measured by the fixed counter-moving crystal method, at a scan rate of $0.5^{\circ} / \mathrm{min}$; the scan range was $1^{\circ}$ for the $0 k l-3 k l$ levels and $2^{\circ}$ for the remaining $4 k l-7 k l$ levels. The backgrounds were measured for 20 seconds at the extreme points of the scan range. During the data collection the intensities of some zero-level reflections were tested periodically to see whether some crystal decay was taking place. However no significant variations of these standards were noticed. The final experimental set, after suppression of all the measurements having a relative statistical error $\sigma(\mathrm{I}) / I>0.25$, was one consisting of 1800 independent intensities. The absorption correction was computed by the Busing and Levy method ${ }^{13}$. Six crystal faces were carefully measured on a Stoe optical goniometer equipped with a Leitz screw micrometer. Because of the relatively low linear absorption coefficient ( $\mu=33.2 \mathrm{~cm}^{-1}$ ) a 216 points sampling $\left(6^{3}\right)$ was adopted; the resulting transmission factors were in the range $0.65-0.72$. The data were reduced to $F_{o}$ values on an arbitrary scale by correcting them for Lorentz and polarization factors. The latter was evaluated by taking into account the partial polarisation of the monochromatized incident beam ${ }^{14}$.

STRUCTURE SOLUTION AND REFINEMENT
The structure has been solved by applying the $\sigma$-two procedure of Karle and Hauptmann ${ }^{15}$ to a set of $220 E$ values greater than 1.8. The detailed procedure used
here has been reported elsewhere ${ }^{6}$. In the present case only one of the 16 signed $E$-sets obtained after convergence of the signing procedure appeared to be promising on the basis of the number of cycles necessary to get to convergence. The corresponding Fourier map did show quite clearly the positions of the six independent metal atoms. A structure factor computation with the contribution of these atoms only gave an $R$ factor of 0.36 . A difference synthesis did then reveal all the light atoms of the structure.

The refinement was done by least squares in the block-diagonal approximation because of the limited computing facilities in the author's laboratory. The blocks were $9 \times 9$ for the heavy atoms treated anisotropically and $4 \times 4$ for the light ones, treated isotropically. The scale factor was correlated with the mean overall temperature factor of the structure in a $2 \times 2$ block $^{16}$. The minimized function was $\Sigma w\left(F_{\mathrm{o}}-k \cdot\left|F_{\mathrm{c}}\right|\right)^{2}$.

The atomic scattering factors used in the least-squares as well as in the structure factors computations, were chosen as follows: for rhodium the T.F.D. statistical model curves ${ }^{17}$ corrected for the real part of the anomalous dispersion, for oxygen, nitrogen and carbon the curves from the self-consistent variational method ${ }^{18}$. The weights $w$ were the reciprocal of the $F_{0}$ 's variance, this variance was computed from the variance of the corresponding $F_{o}^{2}$ assumed to be ${ }^{19} V\left(F_{o}^{2}\right)=V_{\mathrm{cs}}\left(F_{o}^{2}\right)+\left(A \cdot F_{o}^{2}\right)^{2}$, where $V_{c s}\left(F_{0}^{2}\right)$ is the counting statistical variance and $A$ is a parameter chosen in such a way as to render the distribution of the mean weighted quadratic error approximately constant as a function of $\sin \theta / \lambda$ and of $F_{0}$. In this case a value of $A=0.03$ was satisfactory.

The refinement has converged to the following reliability indices:

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R=\frac{\Sigma\left(\left|F_{\mathrm{o}}-k \cdot\right| F_{\mathrm{c}}| |\right)}{\Sigma F_{\mathrm{o}}}=0.050 \text { and } R \cdot w=\left\{\Sigma w\left(F_{\mathrm{o}}-k \cdot\left|F_{\mathrm{c}}\right|\right)^{2} / \Sigma w \cdot F_{\mathrm{o}}^{2}\right\}^{\frac{1}{2}}=0.070
$$

A final difference Fourier has not shown peaks higher than $1 \mathrm{e}^{-} / \AA^{3}$. The $\left|F_{\mathrm{c}}\right|$ and the $F_{\mathrm{o}} / k$ values obtained after convergence are listed in Table 1. The positional and thermal anisotropical parameters of the six rhodium atoms are colfected in Table 2; the positional and thermal isotropic parameters of the $\mathrm{C}, \mathrm{N}$ and O atoms are collected in Table 3. The e.s.d. of each parameter reported in these Tables is derived from the appropriate diagonal element of the L.S. blocks.

## DESCRIPTION OF THE STRUCTURE

One of the metal atoms of the crystallographically independent part of the anion at $1.41 \AA$ from a center of symmetry ( $a$ type position in space group $P 2_{1} / c$ ). This fact implies the existence of a bond twice as long between two corresponding metal atoms of two units $\left[\mathrm{Rh}_{6}(\mathrm{CO})_{15}\right]^{-}$related by inversion. The structure of the anion lies at $1.41 \AA$ from a center of symmetry ( $a$ type position in space group $P 2_{1} / c$ ). through a metal-metal bond. Two CO groups are found in such a position as to form bridges on the two inversion-related metal atoms. Fig. 1 is a perspective view of the anion, a detailed discussion of it is given in the next section of this paper.

The main axis of the anion is almost coincident with $b$, the unique axis of the cell. Further, atoms $\mathrm{Rh}_{1}, \mathrm{Rh}_{2}, \mathrm{Rh}_{3}$ and $\mathrm{Rh}_{4}$ are approximately coplanar with the (100) plane. The crystal packing is formed by anionic layers separated by a distance equal to $9.42 \AA=a \cdot \sin \beta$. The cavities of this packing are occupied by the cations









































[^0]TABLE 1 (continued)

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TABLE 2
final atomic coordinates for anisotropically refined atoms in $\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]_{2}\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]$
All values $\times 10^{4}$; anisotropic temperature factors are in the form: $\exp \left\lceil-\left(h^{2} \cdot \beta_{11}+k^{2} \cdot \beta_{22}+l^{2} \cdot \beta_{33}+h \cdot k \cdot \beta_{12}+h \cdot l \cdot \beta_{13}+k \cdot l \cdot \beta_{23}\right)\right]$.

| Atoms | $x / a$ | $y / b$ | z/c | $\beta_{11}$ | $\beta_{12}$ | $\beta_{13}$ | $\beta_{22}$ | $\beta_{23}$ | $\beta_{3,}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rh1 | 19 (2) | 815 (1) | -68 (1) | 56 (2) | 0 (2) | 4 (1) | 14(0) | -i (0) | 16 (0) |
| Rh 2 | 384 (3) | 2131 (1) | 909 (1) | 199 (3) | -11(2) | 11 (2) | 22 (0) | -9(1) | 19 (0) |
| Rh3 | 41 (2) | 3132 (1) | -315(1) | 84 (3) | 0 (2) | 8 (2) | 16 (0) | -2(1) | 28 (0) |
| Rh4 | -323(2) | 1886(1) | -1298(1) | 79 (2) | -1 (2) | 0 (2) | $19(0)$ | 1 (1) | 18 (0) |
| R15 | -1971 (2) | 2043 (1) | -78(1) | 65 (2) | 10 (2) | 18 (2) | 20 (0) | 2 (1) | 27 (0) |
| Rh6 | 2038 (2) | 1982 (1) | -301 (1) | 64 (2) | -6 (2) | 9 (2) | 19 (0) | -4(1) | 21 (0) |

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TABLE 3
final atomic coordinates for isotropically refineid atoms in $\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]_{2}\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]$ All values $\times 10^{4}$

| Atoms | $x / a$ | $y / b$ | $z / c$ | $\mathrm{B}\left(\AA^{2}\right)$ | Atoms | $x / a$ | $y / b$ | 2/c | $B\left(\AA^{2}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C 11 | -881 (27) | 73 (11) | 588 (12) | 2.3 (4) | C 61 | 3159 (37) | 1362 (17) | -829 (16) | 4.6 (6) |
| O 11 | -1658(21) | 120 (9) | 1052 (9) | 3.7 (3) | 061 | 3959 (29) | 1061 (13) | -1188(13) | 6.7 (5) |
| C 21 | -248 (33) | 1565(14) | 1731 (14) | 3,5(5) | C 62 | 3536 (26) | 2568 (11) | 239 (11) | 1.8 (3) |
| O21 | -585 (26) | 1225(11) | 2222 (11) | 5.3 (4) | O 62 | 4359 (27) | 2956 (11) | 493 (12) | 5.8 (5) |
| C 22 | 1803 (34) | 2660(16) | 1524 (15) | 4.1 (6) | C126 | 1887 (27) | 1092 (11) | 633 (12) | 2.1 (4) |
| O 22 | 2744 (30) | 2917 (13) | 1947 (13) | 6.6 (5) | 0126 | 2741 (22) | 850 (13) | 1068 (10) | 4.3 (4) |
| C 31 | 1201 (39) | 3901 (18) | 68 (17) | 5.1 (7) | C145 | -1759 (28) | 980 (12) | -789 (12) | 2.4 (4) |
| O31 | 2064 (31) | 4371 (14) | 337 (14) | 7.5 (6) | 0145 | -2739 (21) | 676 (9) | -1158(9) | 3.8 (3) |
| C 32 | -1139 (33) | 3801 (14) | -903(15) | 3.6 (5) | C235 | -1170(35) | 2979 (15) | 665 (15) | 3.9 (5) |
| O 32 | -1924 (27) | 4197(12) | -1310(12) | 6.3 (5) | 0235 | -1751 (25) | 3467 (11) | 1035 (11) | 4.8 (4) |
| C 41 | 401 (31) | 1194(13) | - 1974 (14) | 3.1 (5) | C346 | 1287 (32) | 2723 (14) | - 1221 (14) | 3.2 (5) |
| 041 | 699 (26) | $817(11)$ | -2449 (11) | 5.5 (4) | 0346 | 1846 (22) | 3091 (10) | - 1669 (10) | 3,9(3) |
| C 42 | -1819 (32) | 2260(13) | - 2021 (14) | 3.0 (4) | N | 4106 (28) | 50 (12) | 3090 (12) | 3.8 (4) |
| O 42 | -2725 (29) | 2485(13) | -2462(13) | 6.6 (5) | CN1 | 2489 (39) | 48 (16) | 2883 (17) | 5.1 (6) |
| C 51 | -3146 (33) | 1587(14) | 583 (14) | 3.5 (5) | CN2 | 4487 (45) | 257 (20) | 3979 (19) | 6.4 (8) |
| 051 | -3871 (29) | 1348 (13) | 1011(13) | 6.6 (5) | CN3 | 4722 (49) | 667 (21) | 2613 (22) | 7.5 (9) |
| C 52 | -3459 (38) | 2486 (17) | -713(17) | 6.9 (6) | CN4 | 4655 (49) | -752 (21) | 2977 (21) | 7.3 (9) |
| 052 | -4356 (28) | 2769 (12) | -1064(12) | 4.2 (5) |  |  |  |  |  |

$\left[\mathrm{N}\left(\mathrm{CH}_{3}\right)_{4}\right]^{+}$which are found to occur in staggered planes interleaved with the anionic planes (the $x$ coordinates of the cations are 0.41 and 0.59 ).

The cations do not depart significantly from their expected tetrahedral shape. The mean value of the N-C bond lengths in this structure is $1.54(3) \AA$, as compared


Fig. 1. Perspective view of the anion $\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]^{2-}$. The orientation of thrce mutually perpendicular diads of point group $D_{2 h}$ is also shown.


Fig. 2. Relationship between $\mathrm{Rh}_{6}(\mathrm{CO})_{16}$ and $\left[\mathrm{Rh}_{12}(\mathrm{CO})_{30}\right]^{2-}$.
with the commonly accepted ${ }^{20}$ value of $1.48 \AA$. An extensive computation of all possible non-bonding interactions has not revealed significant departures from the expected contacts.

THE TRIACONTACARBONYL DODECARHODATE ANION
The sub-unit $\left[\mathrm{Rh}_{6}(\mathrm{CO})_{15}\right]^{-}$which represents the crystallographic independent part of the anion may be thought of as derived from the neutral species $\mathrm{Rh}_{6}(\mathrm{CO})_{16}$ by subtracting a carbonyl group and adding one electron. This hypothetical reaction

is represented in Fig. 2 where the stereochemical relationship among the two species is shown.

Bond distances and angles of the anion, together with e.s.d.'s computed from the corresponding values of the least squares refinement, are reported in Fig. 3, where the atomic labelling which will be used in the following discussion is given.

In the sub-unit, the five metal atoms, which are not involved in the catenation exhibit exactly the same coordination pattern (referred as "tetragonal antiprismatic" by Dahl ${ }^{7}$ ), as in the neutral parent species. The catenating metal atom does not greatly depart form this scheme but is, in fact, nine-coordinated. The $T_{d}$ symmetry of $\mathrm{Rh}_{6}(\mathrm{CO})_{16}$ is destroyed, in the present anion. The presence of two catenated sub-units, obviously, does not admit any of the $C_{3}$ and $S_{4}$ axes of the $T_{d}$ group. More than that, the $\mathrm{Rh}-\mathrm{Rh}$ distances appear to be widely scattered. However their mean value $2.77 \AA$ compares well with the corresponding value ${ }^{7}$ of $2.78 \AA$ in the neutral cluster.

The deformations of the octahedral moiety may be summarized in this way: (i) the four atoms $\mathrm{Rh}_{2}, \mathrm{Rh}_{4}, \mathrm{Rh}_{5}$ and $\mathrm{Rh}_{6}$ defining the "equatorial" plane, occupy the corners of a rectangle with dimension $2.68(1) \AA \times 2.84(1) \AA$; (ii) the distances of these atoms from the catenating $R h$ are all equivalent $[2.84(1) \AA]$ and so do the corresponding distances from the "apical" $R h_{3}[2.75(1) \AA]$. The difference among the two last groups of distances implies a lower bond order for the first one and is a consequence of the higher coordination number around it. In this range of distances the separation of $2.82(1) \AA$ between the two catenated $\mathrm{Rh}_{1}$ atoms appear not unusual. We may note also that none of the present $\mathrm{Rh}-\mathrm{Rh}$ interactions is a novelty; these distances can be compared with the following values: $2.62 \AA$ in $\mathrm{Rh}_{3}(\pi-\mathrm{Cp})_{3}(\mathrm{CO})_{3}{ }^{21}$, $2.68 \AA$ in $\mathrm{Rh}_{2}(\pi-\mathrm{Cp})_{2}(\mathrm{CO})_{3}{ }^{22}, 2.69 \AA$ in Rh metal ${ }^{20}, 2.72 \AA$ in $\mathrm{Rh}_{3}\left(\mathrm{C}_{20} \mathrm{H}_{21}\right)^{23}$, $2.73 \AA$ in $\mathrm{Rh}_{4}(\mathrm{CO})_{12}{ }^{24}, 2.85 \AA$ in $\mathrm{Ru}_{3}(\mathrm{CO})_{12}{ }^{25}, 2.85-2.95 \AA$ in $\mathrm{Ru}_{6} \mathrm{C}(\mathrm{CO})_{14}(\text { Arene })^{26}$.

The deformations described above for the metal octahedra are in agreement with a point group $D_{2 h}$. Fig. 1 shows three mutually perpendicular diads intersecting at the center of symmetry of the anion. The spatial distribution as well as the bonding distances and angles of all the carbonyl groups follow the same symmetry. So, the eight CO groups linearly bonded to the equatorial $R h_{2}, \mathrm{Rh}_{4}, \mathrm{Rh}_{5}$ and $R h_{6}$ exhibit a narrow range of $\mathrm{Rh}-\mathrm{C}$ and $\mathrm{C}-\mathrm{O}$ distances [mean values $1.88(1) \AA$ and 1.14 (1) $\AA$ respectively], whereas the corresponding distances for the groups bonded to the "apical" $\mathrm{Rh}_{3}$ exhibit a different range $[1.82(2) \AA$ and $1.20(2) \AA]$. A comparison of these distances suggest a more pronounced metal-to-carbon $\pi$-bonding in the apical $\mathrm{Rh}-\mathrm{C}-\mathrm{O}$ interactions.

The face bridging groups show a departure from the threefold symmetry they display in $\mathrm{Rh}_{6}(\mathrm{CO})_{16}$. Here again the spread in metal-carbon and carbon-oxygen distances and also the pattern of the angular distortion is in agreement with the $D_{2 h}$ group. The situation can be described as follows: the inner groups, bridging on the catenating and equatorial atoms, display one shorter $[2.05(3) \AA]$ and two longer [2.31(3) $\AA$ ] values; the reverse is true for the outer groups bridging on the apical and equatorial atoms, one distance being longer $[2.22(2) \AA]$ and two shorter $[2.12(2) \AA]$. The mean $\mathrm{C}-\mathrm{O}$ distance for all these groups is $1.19(3) \AA$.

The group bridged on the $\mathrm{Rh}-\mathrm{Rh}$ bond between the two sub-units lies along one of the diads within the standard deviations of its bonding parameters and has the following bond distances: $\mathrm{Rh}-\mathrm{C} 2.00(2) \AA$ and $\mathrm{C}-\mathrm{O} 1.17(3) \AA$.

The spread of the $\mathrm{Rh}-\mathrm{Rh}$ and $\mathrm{Rh}-\mathrm{C}$ bond lengths, which accompanies the
decrease of symmetry, indicates an overall decrease of the electron delocalisation within the $\left[\mathrm{Rh}_{6}(\mathrm{CO})_{15}\right]^{-}$moiety of the cluster; stated in a different way, the connection of the two octahedral sub-units destroyes the degeneracy of sets of molecular orbitals in an analogous way as the conjugation between aromatic rings does with respect to an aromatic $\pi$-system.

## ACKNOWLEDGEMENTS

The Authors wish to thank V. Scatturin and P. Chini (the latter who also has kindly supplied the crystal samples) for their interest and a number of discussions. The financial support of the Italian Consiglio Nazionale delle Ricerche is gratefully acknowledged.

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